

Volumetrically Scanning the Structure of Stray-fields Above Grain-Oriented Electrical-Steel Using a Variably Angled TMR Sensor

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Abstract—a new versatile scanning hardware based on a *Micromagnetics*[®] *STJ-020* MgO-based tunnelling magneto-resistance (TMR) sensor has been developed to volumetrically scan the thin boundary layer above a given sample. An (x,y) planar scan of the surface of a $7.5\text{ mm} \times 7.5\text{ mm}$ sample of grain-oriented (3% Si) electrical steel is presented. Stray fields normal to the surface between -116 and 272 A/m are measured. At $10\text{ }\mu\text{m}/\text{pixel}$ domain and micro-domain structures are seen. At $5\text{ }\mu\text{m}/\text{pixel}$ the micro-domain structures resolve into clear Lancet domains. The domain images presented have greater qualitative similarity with Kerr effect observations than with Bitter technique results. An (x,z) vertical scan along a 2.35 mm transect reveals the perpendicular extent of the stray fields, with the normal components shown to emanate from the domain bodies and extend approximately $40 - 100\text{ }\mu\text{m}$ from the sample surface. With the aim of investigating how the stray fields close back onto the surface, the (x,z) transect is repeated with the sensor at $5, 10, 15$ and 20 degrees from the vertical. For the first time, the stray fields from surface domains viewed by other techniques in only a planar (x,y) projection have been studied in the (x,z) plane perpendicular to the surface.

Keywords—Magnetic Domains; Grain-oriented Steel; Stray Field; Volumetric Scanning; Tunneling Magneto-Resistance; 3D structure

I. INTRODUCTION

Observation of magnetic domains based upon the Bitter technique [1] highlight the domain boundaries due to the presence of stray magnetic fields between domains. The Kerr effect [2], and observational techniques derived from both such as MOIF [3], reveal the body of domains through their interaction with polarized light. It might be expected that magneto-resistive sensors in a scanning microscope, measuring stray fields, would also highlight the edges of domains where the stray fields are presumed strongest. A three-dimensional scanning magnetic imaging system has been developed based on a *Micromagnetics*[®] *STJ-020* MgO-based tunnelling magneto-resistance (TMR) sensor with a $2\text{ }\mu\text{m} \times 4\text{ }\mu\text{m}$ active area at a distance of $7\text{ }\mu\text{m}$ from the tip edge (Fig. 1(a)).

The TMR sensor is mounted, alongside a *Philtec*[®] *RC20* reflection-compensated fibre-optic displacement sensor for topology tracking, within a custom 3D-printed enclosure

attached to a *Parker Automation* based 3-axis positioning arm with $\pm 1\text{ }\mu\text{m}$ precision (Fig. 1). Similar systems have been used [4] to study geological samples [5] and custom permanent magnetic structures [6], but not for domain imaging at this resolution.

Examples of the three-dimensional study of magnetic fields have involved the use of 3-axis Hall-effect sensors [7] in a 3-axis scanning system [6]. Such sensors must comprise three individual sensors mounted orthogonally, either physically or created on a single substrate [8]. Such complex sensors are limited by fabrication technologies to volumes in excess of 1 mm^3 , or are extremely expensive to create. A Hall-cross size of 700 nm is achieved [8] but with very limited scanning area.

Here, an alternative solution to investigating the structure of stray fields is proposed. The single-axis sensor (Fig. 1(b)) is mounted on a precision goniometer (Fig. 1(c)), by which the angle of the sensor can vary. The geometry of the sensor tip is illustrated in Fig. 2. If the variation in minimum sensor active area to surface distance ($z_0 + \delta z$) can be accounted for, and the coincidence of data and noise levels be confidently managed; then the components of the stray field magnetization vector $H_{(x,z)}$ could be determined.

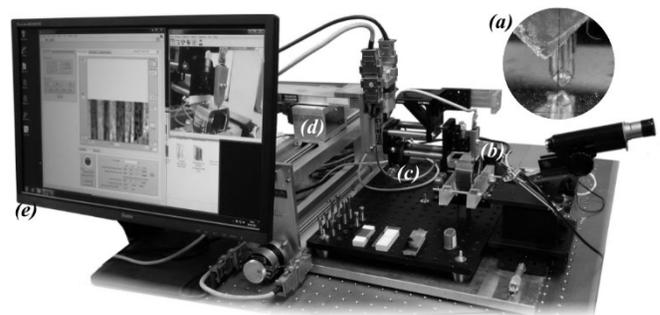


Fig. 1. The system developed is based on a *MicroMagnetics*[®] *STJ-020* tunnelling magneto-resistance sensor with a *Philtec* *RC20* fibre-optic displacement sensor for topology tracking. Figure insert (a) shows the sensor positioned above the reflective sample of electrical steel. The sensors (b) are mounted within a custom 3D-printed enclosure on a precision goniometer (c) attached to a *Parker Automation* based 3-axis positioning arm (d). A comprehensive interface (e) has been developed in *NI LabVIEW*[™] to control the hardware and visualise and interrogate the resulting data.

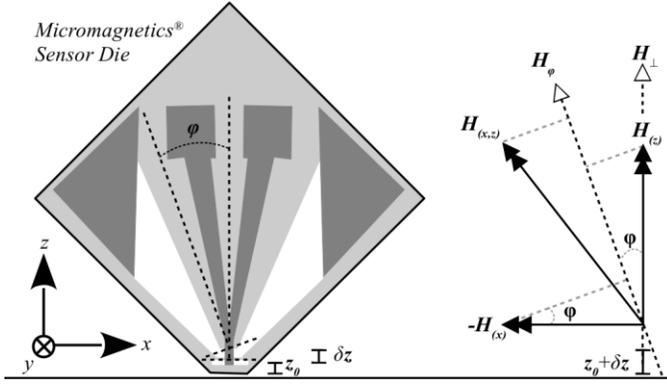


Fig. 2. Illustration of the geometry of the sensor tip. The tip of the sensor die has been polished to bring the active area within $7 \mu\text{m}$ of the tip edge, resulting in a flat tip that enforces an increase in sensor to surface distance ($z_0 + \delta z$) as ϕ increases. $H_{(x)}$ and $H_{(z)}$, the components of the field vector $H_{(x,z)}$ can be derived from H_ϕ , the field measured at angle ϕ and H_\perp , the field measured perpendicular once the difference in field due to δz is determined.

From two measurements of the perpendicular field H_\perp and the field H_ϕ at angle ϕ then

$$H_\phi = H_{(z)} \cos(\phi) - H_{(x)} \sin(\phi) \quad (1)$$

where $H_{(z)}$ and $H_{(x)}$ are the components of $H_{(x,z)}$, and thus

$$H_{(x)} = (H_{(z)} \cos(\phi) - H_\phi) / \sin(\phi) \quad (2)$$

and

$$H_{(z)} = H_\perp - H_{\delta z} \quad (3)$$

with $H_{\delta z}$ being the difference in stray field strength near the surface and at the height $z_0 + \delta z$ above the surface.

II. PLANAR (X,Y) SCAN OF GRAIN-ORIENTED (3% Si) ELECTRICAL STEEL

Planar (x,y) stray field scans of a polished uncoated sample of grain-oriented electrical steel (3% Si B32L) were made, with the sensor perpendicular to the sample surface (angle of sensor to vertical $\phi = 0^\circ$). Stray fields of between -116.0 and 272.0 A/m were recorded normal to the sample surface, where a positive value indicates fields out from the surface, and negative into the surface. The resulting images are presented in Fig. 3 and Fig. 4. Grey (50%) represents either true null-field or a field perpendicular to the sensor axis.

Each pixel in the resultant scans represents the mean of 20 samples of the field through the sensor's active area. Investigation was made into the best compromise between effective signal-to-noise reduction and increased scan time.

Main bar domains are clearly visible with evidence of internal Lancet [9] domains present. The domain images have greater qualitative similarity with Kerr effect (surface domain magnetization) observations than with Bitter technique (domain boundary) observations. Although scanned at a spatial resolution of $5 \mu\text{m}/\text{pixel}$, the system currently has limited resolving power for features smaller than $40 - 50 \mu\text{m}$.

Three factors limit the system resolution; the scanner arm precision (calibrated to $\pm 1 \mu\text{m}$), the width of the sensor's active area ($4 \mu\text{m}$) and the sensor to sample-surface distance (z_0). The sensor tip has been polished to bring the active area within $7 \mu\text{m}$ of the tip edge. With the sensor vertical ($\phi = 0^\circ$), z_0 is still $9 \pm 1 \mu\text{m}$ making this the dominant factor effecting the resolution. Continued efforts to control $z_0 + \delta z$ further include better topology tracking using the *Philtec® RC20* displacement sensor to directly monitor tip height.

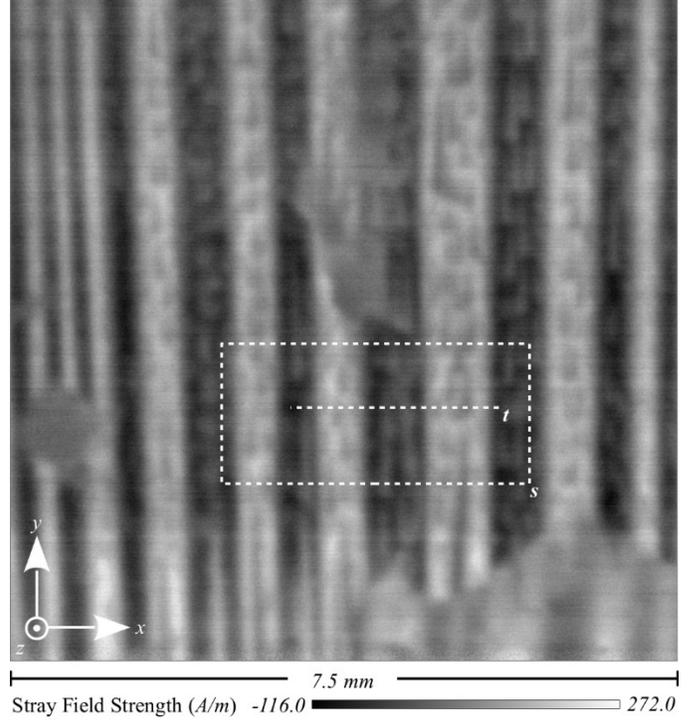


Fig. 3. Planar (x,y) stray field scan of $7.5 \text{ mm} \times 7.5 \text{ mm}$ region of grain-oriented electrical steel (3% Si B32L) at a resolution of $10 \mu\text{m}/\text{pixel}$ with the sensor axis vertical ($\phi=0^\circ$). Stray fields of approximately ± 200 A/m are recorded normal to the sample surface suggesting internal Lancet domains within the clear bar domains. Area "s" indicates the location of the sub-region (Fig. 4), and the line "t" indicates the location of the (x,z) transects (Fig. 6).

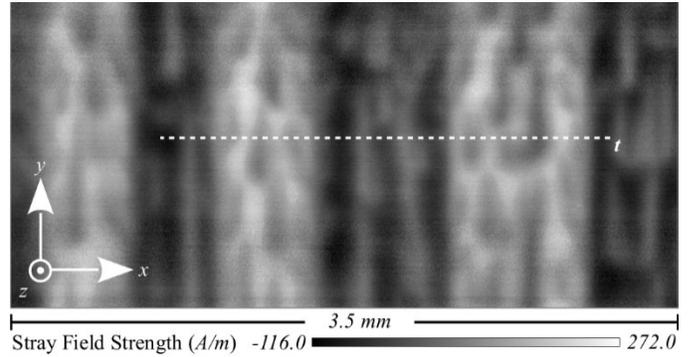


Fig. 4. Planar (x,y) stray field scan of the $3.5 \text{ mm} \times 2.14 \text{ mm}$ sub-region of grain-oriented electrical steel (3% Si B32L) indicated in Fig. 3 at a resolution of $5 \mu\text{m}/\text{pixel}$. Stray fields of approximately ± 200 A/m are recorded normal to the sample surface clearly indicating internal Lancet domains within the main bar domains. The location of the (x,z) transects (Fig. 6) is marked with line "t".

III. TRANSECT (x,z) SCANS OF GRAIN-ORIENTED (3% Si) ELECTRICAL STEEL

Scans in the vertical (x,z) plane have been made along the 2.35 mm transect indicated in Fig. 3 and Fig. 4, to a height of $z_0 + 0.1$ mm.

A. Vertical (x,z) Transect Scan

An illustration of the three-dimensional relationship between the vertical (x,z) scan and the planar (x,y) scan is presented in Fig. 5. The fields directly perpendicular to the sample surface are presented in the first ($\varphi = 0^\circ$) transect of Fig. 6. It can be seen that the fields emanating from the body of the smaller Lancet domains are limited in extent to approximately 40 μm above the surface, before their effect is subsumed by the larger surrounding bar domains. The bar domains remain distinct to greater than 100 μm .

B. Angled (x,z) Transect Scans

Four further scans of the 2.35 mm \times 0.1 mm transect, with a sensor axis angle φ of 5° , 10° , 15° and 20° to the vertical, are also presented in Fig. 6. The intention is to determine the shape of the fields emanating from the Lancet domains and indicate how they may close back to the surface of the sample. At present, because it is difficult to distinguish between zero field or field perpendicular to the sensor axis the, results presented in Fig 6 have two possible causes; a rotation of the field vector compared with the sensor axis, or an x-axis positional drift (less than 1 μm) coupled with a decrease in field magnitude caused by the enforced increase δz (Fig. 2) due to the increase in φ .

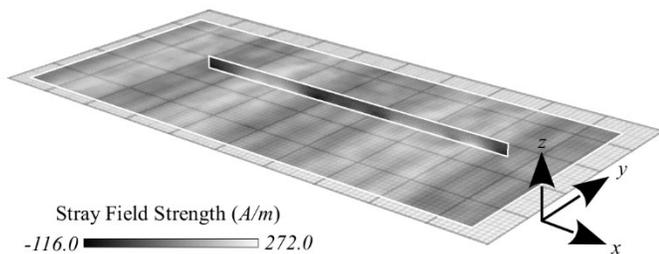


Fig. 5. Illustration of the three-dimensional relationship between the vertical (x,z) scans and the planar (x,y) scan. The (x,y) plane is presented at 50% opacity to emphasise the (x,z) plane.

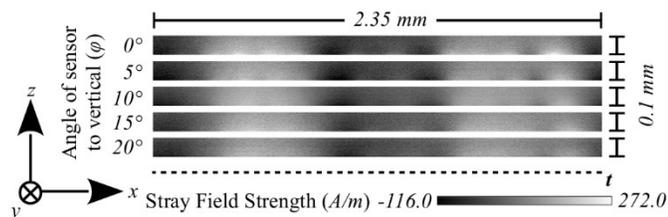


Fig. 6. Five (x,z) scans of the same 2.35 mm \times 0.1 mm transect, indicated by “t” in Fig. 4, with sensor axis angle φ of between 0° and 20° to the vertical. The transects have an x-axis resolution of $5\mu\text{m}/\text{pixel}$ and z-axis resolution of $2\mu\text{m}/\text{pixel}$ with 5:2 ratio non-square pixels to maintain the correct relative dimensions. The resulting drift of the field patterns have two possible causes, discussed in the text, which are currently difficult to isolate.

IV. CONCLUSIONS

A sample of grain-oriented electrical steel (3% Si) has been successfully scanned, indicating Lancet domain structures of width 100 – 150 μm within bar domains of width 600 – 800 μm . The fields emanating from the body of these domains are shown to be ± 200 A/m and extend 40 – 100 μm above the surface of the sample.

With the minimum sensor active area to surface distance z_0 of 9 ± 1 μm , and the fields of some of the smaller Lancet domain fields shown to extend no more than $z_0 + 40$ μm from the surface of the sample (Fig. 6), this distance represents a failure to resolve the lowest 16% of the field. When φ is increased, δz is enforced further by the flat edge of the tip.

Levels of noise remain in the data which, while qualitatively manageable by eye, cause difficulties for stricter quantitative geometric calculations. More work is needed to manage noise levels and the error margins in coincidence of repeated scans, before the components of the field at each point can be confidently isolated.

However, for the first time, the stray fields from surface domains viewed by other techniques [1, 2, 3] in only a planar (x,y) projection have been studied in the (x,z) plane perpendicular to the surface. The strength and extent of the fields have been quantified and progress made towards demonstrating their shape, and how it varies spatially in the surface boundary layer with the underlying domain structure.

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